

EE 213, Microscopic Nanocharacterization of Materials

Lecture 5. W2016

Mike Isaacson, Baskin 237

Email: msi@soe.ucsc.edu

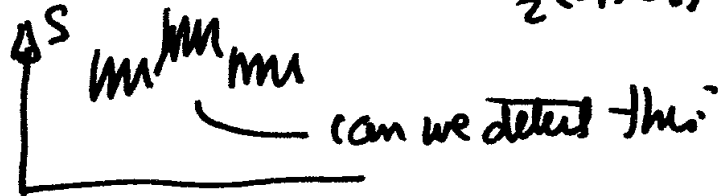
Tele: 831-459-3190

Admin. Asst. Rachel Cordero: rcordero@soe.ucsc.edu, 831-459-2921

Contrast of signals

ASIDE/ contrast detectable using BSE
either for Z contrast or threshold contrast

generally define $C = \frac{S_1 - S_2}{\frac{1}{2}(S_1 + S_2)} = \frac{\text{diff}}{\text{AVG}}$



we usually say $\frac{\Delta S}{S} > \frac{1}{\sqrt{S}}$ — assume poisson counting noise

$$\text{ie } C > \frac{1}{\sqrt{S/N}} \quad || \quad \frac{S}{N} \sim \sqrt{S}$$

min contrast defined as

$$C_{\min} = \frac{k}{(S/N)} \quad \text{unit } k \text{ between 1 and 5}$$

see Box: Vision: Human & Electronic, 1974 (came out in 1948)
early days of TV

Albert Rose

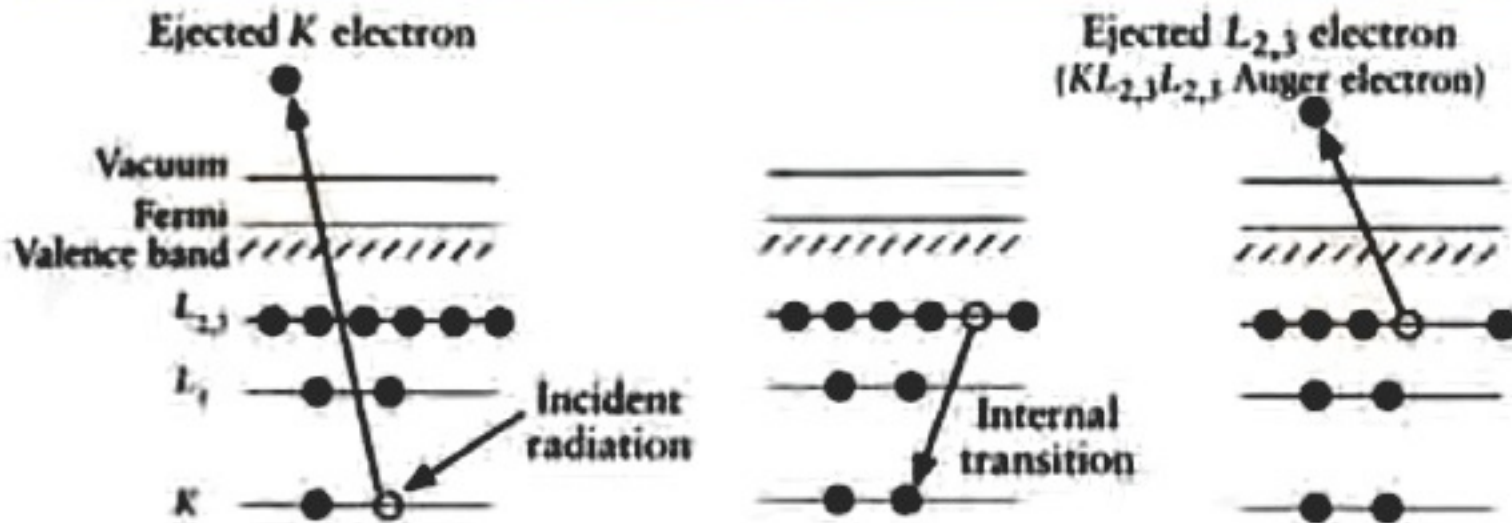
VISION

HUMAN AND ELECTRONIC



Auger Spectroscopy Nomenclature

Auger Electron Spectroscopy



First approximation: $E_{abc} = E_a - E_b - E_c$

Nucleus in ionized atom has one more + charge and looks like atom with Z+1

Better approximation:

$$E_{abc}(Z) = E_a(Z) - 0.5[E_b(Z) + E_b(Z+1)] - 0.5[E_c(Z) + E_c(Z+1)]$$

Atomic Electron Binding Energies.

K. Siegbahn, et.al. (1967). "Atomic, Molecular and Solid State Structure Studied by Means of Electron Spectroscopy. (Almqvist and Wiksells, Uppsala)

2. ELECTRON BINDING ENERGIES FOR AES

(eV)

	1s _N	2s _N	2p _N	2p _N	3s _N	3p _N	3p _N	3d _N	3d _N	4s _N	4p _N	4p _N	4d _N	4d _N	4d _N	5s _N	5p _N	5p _N	5d _N	5d _N
	K	L ₁	L ₂	L ₃	M ₁	M ₂	M ₃	M ₄	M ₅	N ₁	N ₂	N ₃	N ₄	N ₅	N ₆	O ₁	O ₂	O ₃	O ₄	O ₅
1H	13																			
2He	21																			
3Li	55																			
4Be	111																			
5B	188			5																
6C	284			7																
7N	399			9																
8O	532	24		7																
9F	686	31		9																
10Ne	867	45		18																
11Na	1 072	63		21	1															
12Mg	1 205	89		52	2															
13Al	1 560	118	74	73	1															
14Si	1 839	149	100	99	8		3													
15P	2 149	189	136	135	16		10													
16S	2 472	229	185	184	16		8													
17Cl	2 823	270	202	200	18		7													
18Ar	3 203	320	247	245	25		12													
19K	3 608	377	297	294	34		18													
20Ca	4 038	438	350	347	44		26													
21Sc	4 493	500	407	402	54		32													
22Ti	4 965	564	461	455	59		34													
23V	5 465	628	520	513	66		38													
24Cr	5 989	695	584	575	74		43													
25Mn	6 539	769	652	641	84		49													
26Fe	7 114	846	723	710	95		56													
27Co	7 709	926	794	779	101		60													
28Ni	8 333	1 008	872	855	112		68													
29Cu	8 979	1 096	951	931	120		74													
30Zn	9 659	1 191	1 044	1 021	137		87													
31Ga	10 367	1 296	1 143	1 116	156		107													
32Ge	11 104	1 413	1 249	1 217	181		129													
33As	11 867	1 547	1 359	1 323	204		147													
34Se	12 658	1 694	1 476	1 436	232		168													
35Br	13 474	1 852	1 596	1 550	257		189													
36Kr	14 326	1 921	1 727	1 675	289		223													
37Rb	15 200	2 065	1 864	1 805	322		248													
38Sr	16 105	2 216	2 007	1 940	358		280													
39Y	17 039	2 373	2 155	2 000	395		313													
40Zr	17 998	2 532	2 307	2 223	431		345													
41Nb	18 984	2 698	2 465	2 371	469		379													
42Mo	20 000	2 866	2 625	2 520	505		410													
43Tc	21 044	3 042	2 793	2 672	544		445													
44Ru	22 117	3 224	2 967	2 838	585		483													
45Rh	23 220	3 412	3 146	3 004	627		521													

	4s _N	4p _N	4d _N	4d _N	4d _N
	P ₁	P ₂	P ₃	P ₄	P ₅
43Tc	3		1		
43Tc	8		3		
44Ru	12		5		
45Rh	18		8		
46Pd	26		11		
47Ag	34		15		
48Cd	44		19		
50Tb	60	09	13	2	2

Atomic Electron Binding Energies.

K. Siegbahn, et.al. (1967). "Atomic, Molecular and Solid State Structure Studied by Means of Electron Spectroscopy. (Almqvist and Wiksells, Uppsala)

	(eV)																								
	1s _N	2s _N	2p _N	2p _N	3s _N	3p _N	3p _N	3d _N	3d _N	4s _N	4p _N	4p _N	4d _N	4d _N	4f _N	4f _N	4f _N	5s _N	5p _N	5p _N	5d _N	5d _N	5d _N		
	K	L ₁	L ₂	L ₃	M ₁	M ₂	M ₃	M ₄	M ₅	N ₁	N ₂	N ₃	N ₄	N ₅	N ₆	N ₇	O ₁	O ₂	O ₃	O ₄	O ₅	O ₆	O ₇		
46 Pd	24 350	3 405	3 331	3 173	670	559	531	340	335	86	51				1										
47 Ag	25 514	3 096	3 524	3 351	717	602	571	373	367	95	62	56			3										
48 Cd	26 711	4 013	3 737	3 526	770	651	617	411	404	104		67			9						2				
49 In	27 940	4 228	3 938	3 730	826	702	664	451	443	122		77			16						1				
50 Sn	29 200	4 465	4 156	3 929	884	757	715	494	485	137		89			24						1				
51 Sb	30 491	4 699	4 381	4 132	944	812	766	537	528	152		99			32						1				
52 Te	31 814	4 928	4 612	4 341	1 006	870	819	582	572	164		110			40						2				
53 I	33 170	5 188	4 852	4 557	1 072	931	875	631	620	186		123			50						3				
54 Xe	34 561	5 452	5 104	4 782	1 145	999	937	685	672	204		147			62						7				
55 Cs	35 985	5 712	5 360	5 013	1 217	1 065	998	740	726	231	172		162	79							23	13		12	
56 Ba	37 441	5 987	5 624	5 247	1 293	1 137	1 067	796	781	253	192		180	93							48	17		15	
57 La	38 925	6 267	5 891	5 483	1 362	1 205	1 124	849	832	271	206		192		99						33		15		
58 Ce	40 444	6 549	6 165	5 724	1 435	1 273	1 186	902	884	290	224		204		111						38		20		
59 Pr	41 991	6 835	6 441	5 965	1 511	1 338	1 243	951	931	305	237		218		114						38		23		
60 Nd	43 569	7 126	6 772	6 208	1 576	1 403	1 298	1 000	978	316	244		225		118						38		22		
61 Pm	45 185	7 428	7 013	6 460	1 650	1 472	1 357	1 052	1 017	331	255		237		121						38		22		
62 Sm	46 835	7 737	7 312	6 717	1 724	1 542	1 421	1 107	1 011	347	267		249		130						29		22		
63 Eu	48 519	8 052	7 618	6 977	1 800	1 614	1 481	1 161	1 131	360	284		257		134						32		22		
64 Gd	50 239	8 376	7 931	7 243	1 881	1 689	1 544	1 218	1 164	376	289		271		141						36		21		
65 Tb	51 996	8 708	8 252	7 515	1 968	1 768	1 612	1 276	1 242	390	311		286		148						40		26		
66 Dy	53 788	9 047	8 581	7 790	2 047	1 842	1 676	1 332	1 295	416	332		293		154						63		26		
67 Ho	55 618	9 395	8 919	8 071	2 128		1 741	1 391	1 251	436	343		304		161						51		20		
68 Er	57 486	9 772	9 265	8 207	2 207	2 005	1 812	1 453	1 409	449	366		320	177							60		29		
69 Tm	59 390	10 116	9 618	8 468	2 297	2 090	1 885	1 515	1 468	472	386		337		180						53		32		
70 Yb	61 332	10 488	9 978	8 843	2 397	2 172	1 949	1 576	1 527	487	386		343	197							53		23		
71 Lu	63 314	10 870	10 349	9 244	2 491	2 264	2 024	1 640	1 589	506	410		359	205							57		28	5	
72 Hf	65 351	11 272	10 739	9 561	2 601	2 365	2 108	1 716	1 642	538	437		380	224							65	38		7	
73 Ta	67 417	11 680	11 136	9 881	2 708	2 469	2 194	1 793	1 725	566	465		405	242							71	45		6	
74 W	69 525	12 099	11 542	10 205	2 820	2 575	2 281	1 872	1 800	595	492		426	259							73	47		6	
75 Re	71 677	12 537	11 957	10 525	2 932	2 682	2 367	1 949	1 883	625	518		445	274							83	46		4	
76 Os	73 871	12 988	12 385	10 871	3 049	2 792	2 458	2 031	1 960	655	547		469	290							50	84		0	
77 Ir	76 111	13 419	12 824	11 215	3 174	2 899	2 551	2 116	2 044	690	577		495	312							60	96		4	
78 Pt	78 395	13 880	13 273	11 564	3 298	3 027	2 646	2 202	2 131	724	608		519	331							70	102		2	
79 Au	80 725	14 353	13 733	11 918	3 425	3 150	2 743	2 291	2 206	759	644		546	352							83	108		3	
80 Hg	83 103	14 829	14 209	12 284	3 562	3 279	2 847	2 385	2 295	800	677		571	379							99	120		7	
81 Tl	85 531	15 347	14 698	12 657	3 704	3 416	2 957	2 485	2 380	846	722		609	407							118	137	100		13
82 Pb	88 005	15 841	15 200	13 025	3 851	3 554	3 067	2 586	2 484	894	764		645	435							138	148	105		20
83 Bi	90 526	16 388	15 708	13 418	3 999	3 697	3 177	2 688	2 580	938	806		679	464							160	117			25
84 Po	93 105	16 929	16 244	13 814	4 149	3 854	3 302	2 798	2 683	995	851		705	500							177	132			31
85 At	95 730	17 492	16 785	14 214	4 312	4 008	3 426	2 909	2 787	1 042	886		740	533							196	148			40
86 Rn	98 404	18 049	17 327	14 619	4 472	4 159	3 538	3 022	2 892	1 097	929		768	567							214	164			48
87 Fr	101 137	18 629	17 906	15 021	4 652	4 327	3 662	3 136	3 000	1 153	980		810	603							234	182			58
88 Ra	103 922	19 237	18 484	15 444	4 822	4 490	3 792	3 248	3 105	1 208	1 058		879	636							254	200			68
89 Ac	106 755	19 840	19 083	15 871	5 002	4 656	3 909	3 370	3 213	1 269	1 080		900	675							272	215			80
90 Th	109 651	20 472	19 693	16 300	5 182	4 831	4 046	3 491	3 322	1 320	1 168		968	714							290	229			88

Estimating Auger Electron Energies

1st approx: $E_{ABC}(z) = (E_A(z) - E_B(z)) - E_C(z)$

initial ionization → shell of excited electron.
 shell where e⁻ came from to fill hole

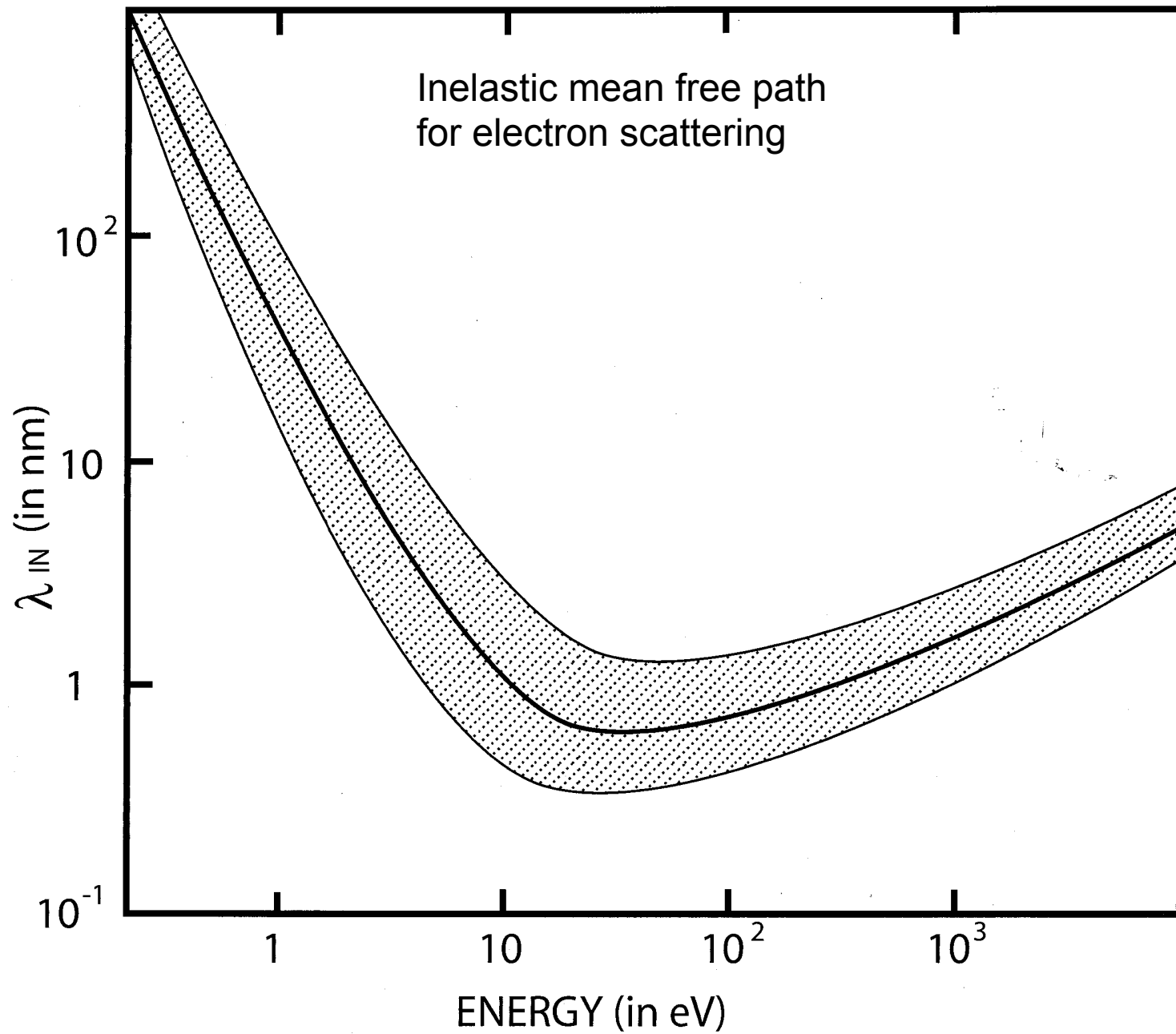
better approx: $E_{ABC}(z) = E_A(z) - \frac{1}{2} [E_B(z) + E_B(z+1)]$
 $-\frac{1}{2} [E_C(z) + E_C(z+1)]$

example: KL_2L_3 transition in Al ($z=13$)

shell	Al ($z=13$)	Si ($z=14$)
E_K	1560 eV	1839 eV
E_{L_1}	118 eV	149 eV
E_{L_2}	74 eV	100 eV
E_{L_3}	73 eV	99 eV

\therefore 1st approx / $E_{KL_2L_3}(Al) = 1560 - 74 - 73 = 1413 \text{ eV}$ // generally higher

better / $E_{KL_2L_3}(Al) = E_K(13) - \frac{1}{2} [E_{L_2}(13) + E_{L_2}(14)]$
 $-\frac{1}{2} [E_{L_2}(13) + E_{L_3}(14)]$
 $= 1560 - \frac{1}{2} [74 + 100] - \frac{1}{2} [73 + 99]$
 $= \underline{\underline{1387 \text{ eV}}}$ // ← actual \approx < 1% diff.



From Seah and Dench, 1979. Surf. and Interface Anal.1.36

① Auger Electron Analysis

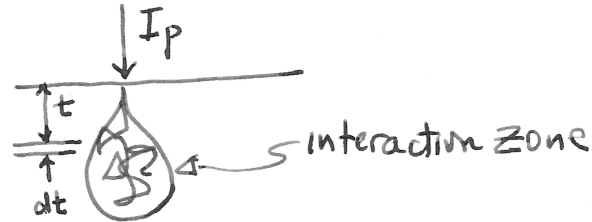
read:

M. Isaacson,

Ultramicroscopy, 4(1993), 171-178

$$S = NJ\sigma YF$$

genl nihar quant. eqn.



$$J = I_p / \left(\frac{\pi}{4} d_B^2 \right) \quad \text{current density striking sample /}$$

Auger electrons that escape

$$\frac{I_{\text{esc}}}{I_{\text{produced}}} = e^{-t/\lambda_{\text{esc}}}$$

note:

$$\lambda_{\text{esc}} \approx \lambda_n$$

but not exactly

∴ # produced in thickness dt at depth t which "escape" solid

$$dN(t) = n(t) A dt \quad \text{area at depth } t$$

$$\therefore NJ = \int_0^{\infty} n(t) I_p e^{-t/\lambda} dt \quad \text{NOTE // } I_p \text{ vs } I_B$$

2- Auger Electron Analysis (cont)

NOTE / probe used that can produce Auger electrons $> I_p$

since BSE can produce Auger also

$$\therefore I_{EFF} = I_p + I_B$$

current in
inc. probe

current due to
BSE electron

how to estimate I_B ?

simple // $I_B \cong I_p [1 + R^*]$

the BSE for particular
"energy loss".

R^* not exactly η , BSE yield
it depends on binding energy

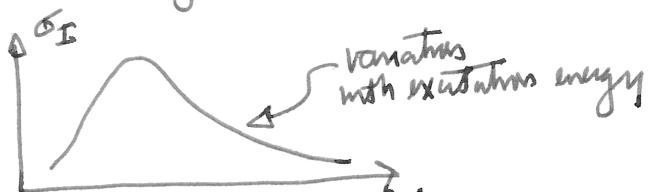
i.e. η of BSE yield depends on energy lost by BSE
before they produce anything. (general 1 shell)

3. Auger Electron Analysis. (cont)

$$S = NJ\sigma YF$$

$$\therefore NJ = \int_0^{\infty} dt n(t) I_p [1 + R^*(t)] e^{-t/\lambda}$$

remembers, the cross section for producing an inner shell ionization looks like



$$\therefore NJ\sigma = \int_{E_w}^{E_p} dE \int_0^{\infty} dt n(t) I_p [1 + R^*(t)] e^{-t/\lambda} \frac{d\sigma(E_w, E_p)}{dE}$$

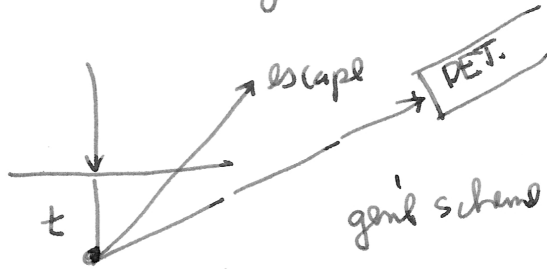
\downarrow energy of w^{th} shell binding energy

this is total signal produced
 now we need to calculate the yield, Y , and
 the fraction that gets collected
 i.e. we need YF

4. Auger Electron Analysis (cont)

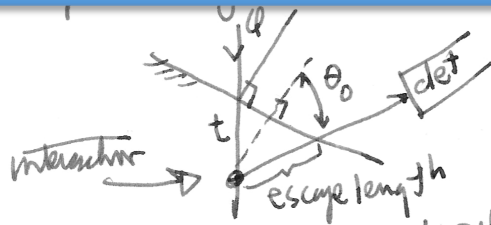
lets now look at F , the efficiency factor ("fudge")

ϕ



$$\frac{I_{esc}}{I_{prod}} = e^{-t/\lambda}$$

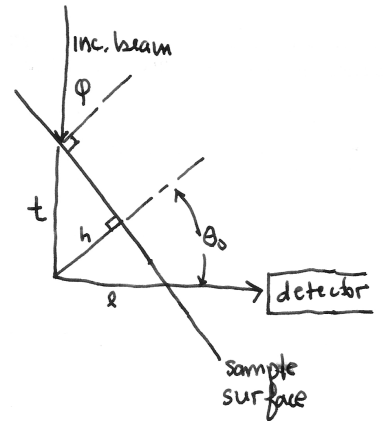
Sample surface normal at ϕ to incident beam



$$I_{esc}/I_{prod} = e^{-\frac{t \cos \phi}{\lambda \cos \theta_0}}$$

$$\equiv F_R e^{-\frac{t \cos \phi}{\lambda \cos \theta_0}}$$

takes into account surface roughness



$$\frac{h}{t} = \cos \phi$$

$$\frac{h}{l} = \cos \theta_0 \implies l = \frac{t \cos \phi}{\cos \theta_0}$$

5. Auger Electron Analysis (cont)

$\therefore S = N I \sigma Y F$ micron eqn.

$$S = \frac{1}{4\pi} \int d\Omega \int_{E_w}^{E_p} dE Y_A \int_0^{\infty} dt n_p(t) \left[1 + R_B^* \right] \frac{d\sigma}{dE} r_B e^{-t/\lambda}$$

integrate over
solid angle
of detector

yield of Auger electrons
we are counting from
A in material.

// $Y_A = \text{Auger Yield}$ //

So, how do we get to our simple equation?

$$S = \frac{n I_p Y_A}{4\pi} \int d\Omega \int_{E_w}^{E_p} dE \int_0^{\infty} dt \left[1 + R_B^* \right] \frac{d\sigma}{dE} r_B e^{-t/\lambda}$$

doesn't depend
on material

[note: we must
assume n doesn't
change thru volume]

$Y_A = \text{Auger yield} = \gamma_w \times \gamma$

average R_B^* and assume $\frac{d\sigma}{dE}$ doesn't change much
with depth from which Auger electrons emerge.

then $\int_0^{\infty} dt e^{-t/\lambda} = \lambda$

5. Auger Electron Analysis (cont)

$$\therefore S_A = \frac{N I_p Y_A}{4\pi} (1 + R_B^*) \underbrace{\sigma(E_i, E_p)}_{\substack{\text{inel x section fn} \\ \text{w shell ionization by inc. electrons} \\ \text{of energy } E_p}} \int d\Omega \Lambda \Gamma_B^*$$

complex factor

$$\text{let } T \Gamma_B^* = \frac{1}{4\pi} \int_{\text{DET}} d\Omega R_R^*$$

transmission function of detector \leftarrow hardest to "estimate"

$$\text{then } S_A = \underbrace{N I_p Y_A}_{NJ} (1 + R_B^*) \underbrace{\sigma(E_i, E_p)}_{\sigma} \underbrace{Y[\Gamma_B^* T]}_F$$

$$\text{then } \frac{S_A}{S_B} = \frac{N_A [stuff]_A}{N_B [stuff]_B}$$

on quantitative eqn //

note / $R_R^* = R_B^*$ tabulated //

NOTE: \downarrow
OA
 \uparrow
OA

still the same
if how to distinguish
the 2 cases.

A near inf. - fms
A far below //

?

Angular quantities

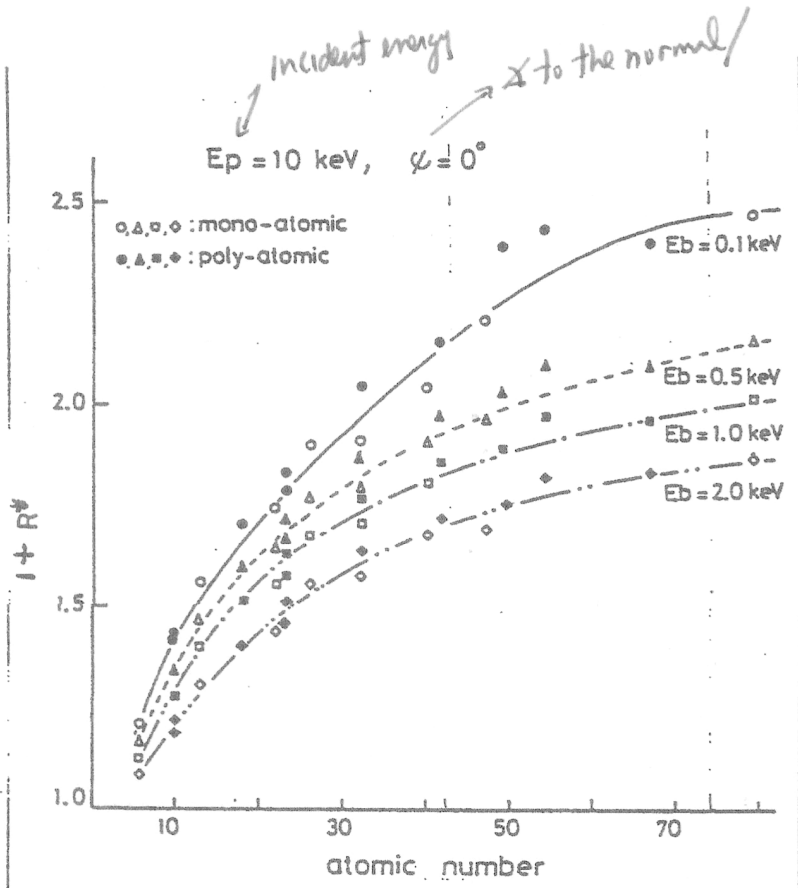
$$S_A = n_A I_p \lambda_A (h R_A^{\frac{1}{2}}) \delta(E_w, E_p) Y[\Gamma_A^{\frac{1}{2}} T]$$
$$= n_A [\text{stuff}]_A I_p$$

$$\text{where } [\text{stuff}]_A = \lambda_A (h R_A^{\frac{1}{2}}) \delta(E_w, E_p) Y[\Gamma_A^{\frac{1}{2}} T]$$

if measurements simultaneous then $I_{PA} = I_{PB} = I_p$ /

$$\therefore \frac{S_A}{S_B} = \frac{n_A [\text{stuff}]_A}{n_B [\text{stuff}]_B}$$

the basis of quantities /

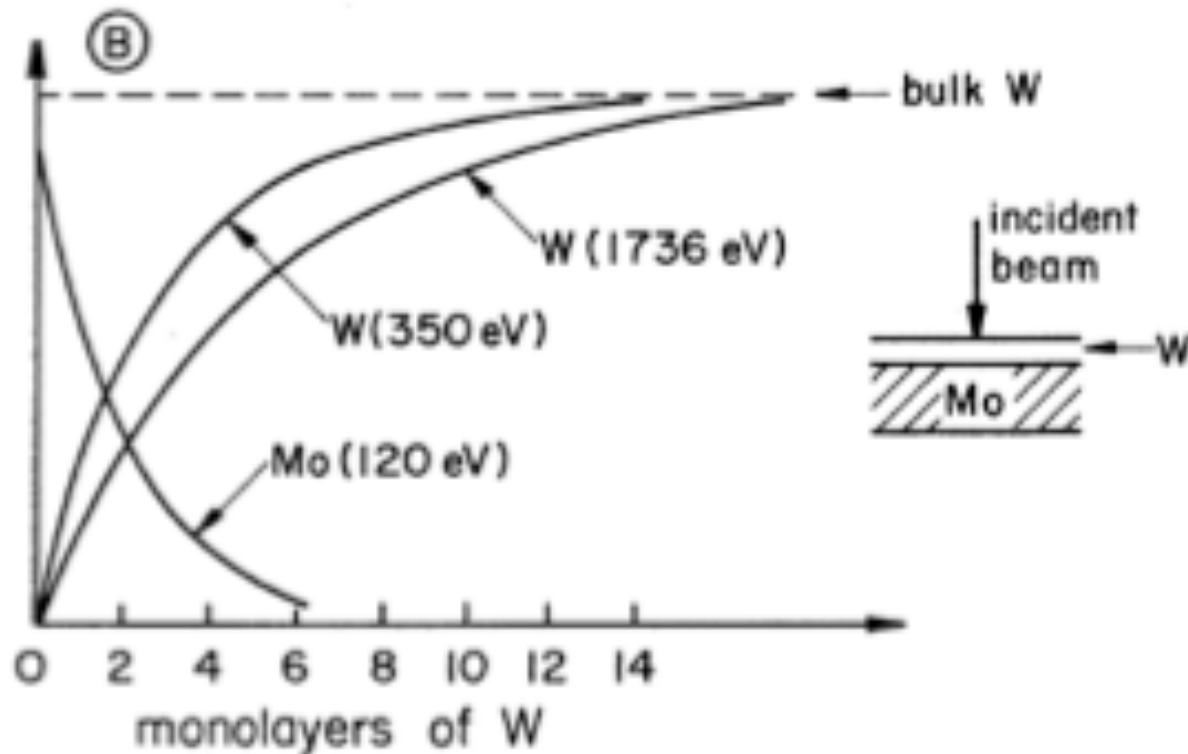


the E_b are the binding energy of the shell from which the original ionization occurred

from Shuniga
Toyota Foundation
Report # I-006 (F83)

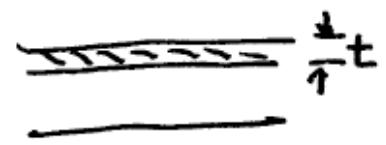
Using AES to Measure Film Thickness

After Tarnag and Wehner. J.Appl.Phys.44 (1973).1534

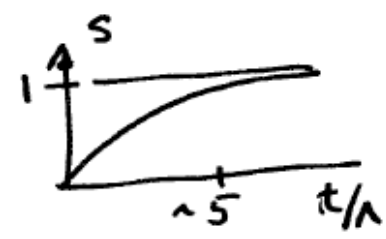


overlayers // Tamy? Wehmen JAP. 44, 1534 (1973)

Mo over W or W over Mo



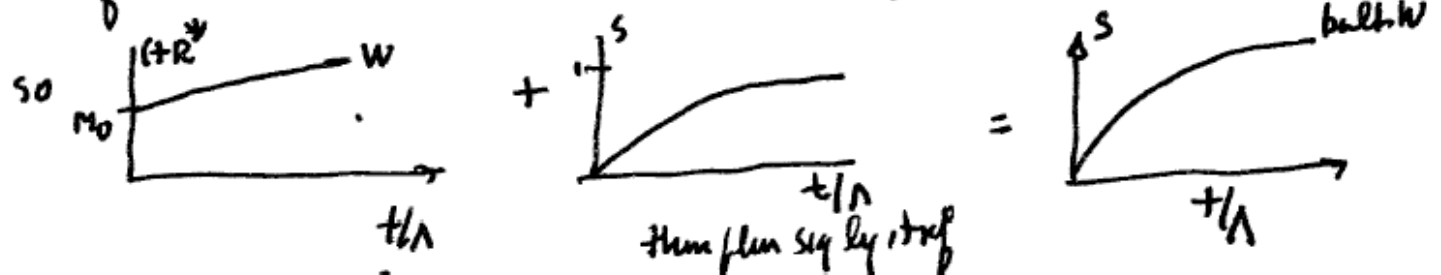
expect signal from them by itself.



$$\int_0^t dt e^{-t/\lambda} = \lambda(1 - e^{-t/\lambda})$$

for thick overlayers $t/\lambda \gg \lambda$ approaches bulk value. - additional Auger at start: $t/\lambda \ll \lambda$ due to substrate, at end due to overlayers.

so if W over Mo // $t/\lambda_W > t/\lambda_{Mo}$ in bulk

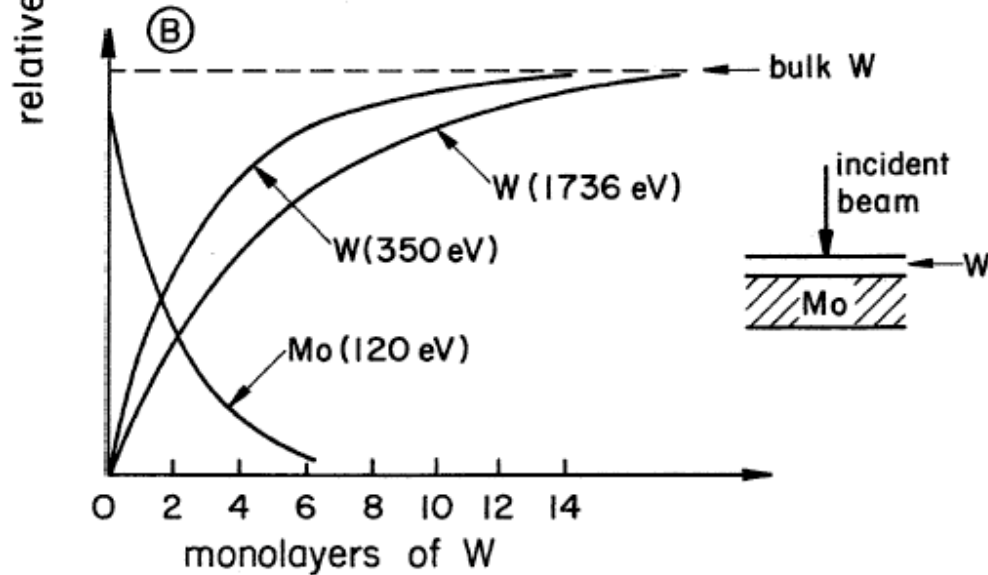
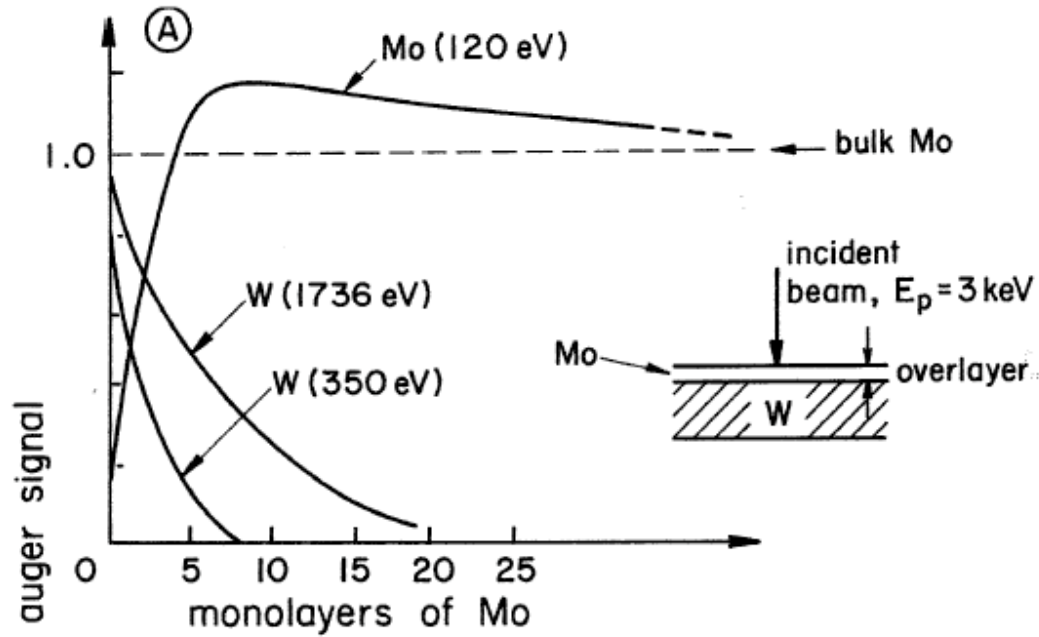


and signal from Mo goes down as overlayers increase

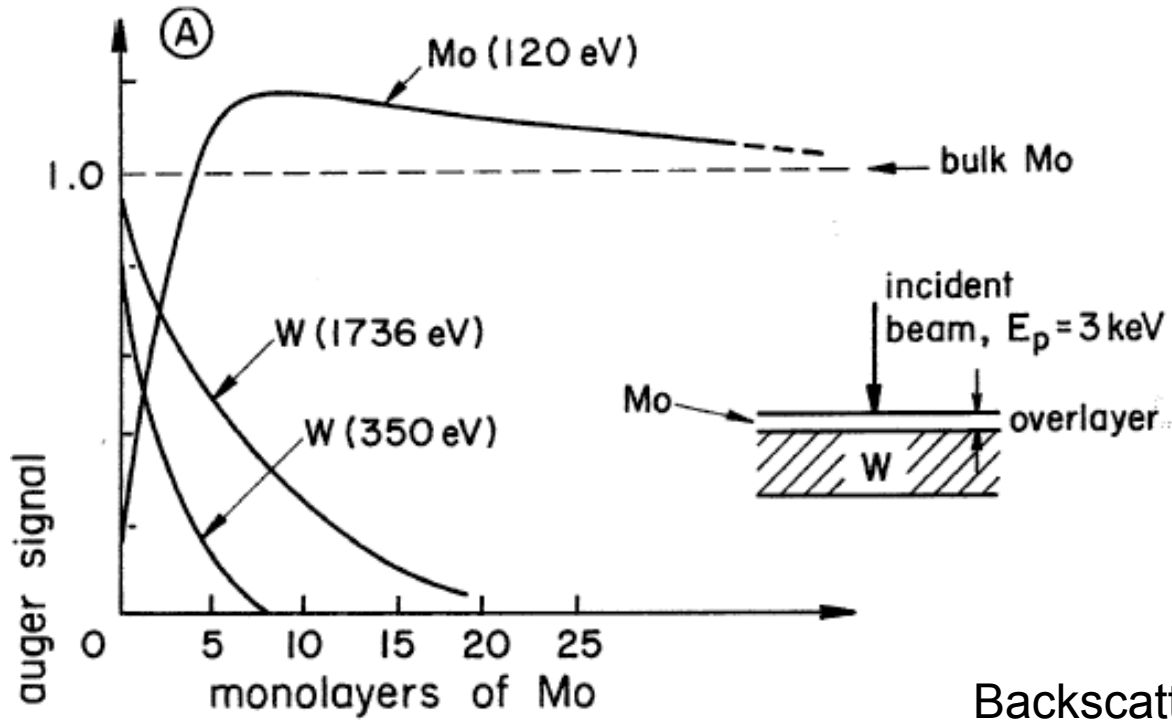


Using AES to Measure Film Thickness

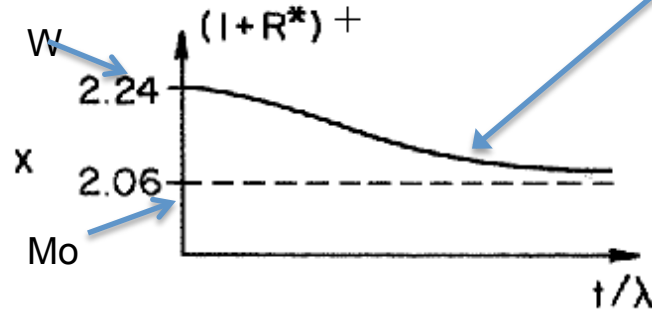
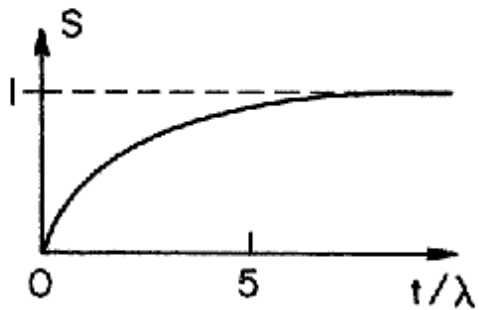
After Tarrg and Wehner. J.Appl.Phys.44 (1973).1534



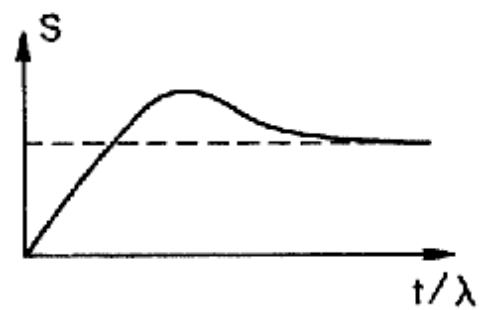
Using AES to Measure Film Thickness (cont)



Backscattering factor

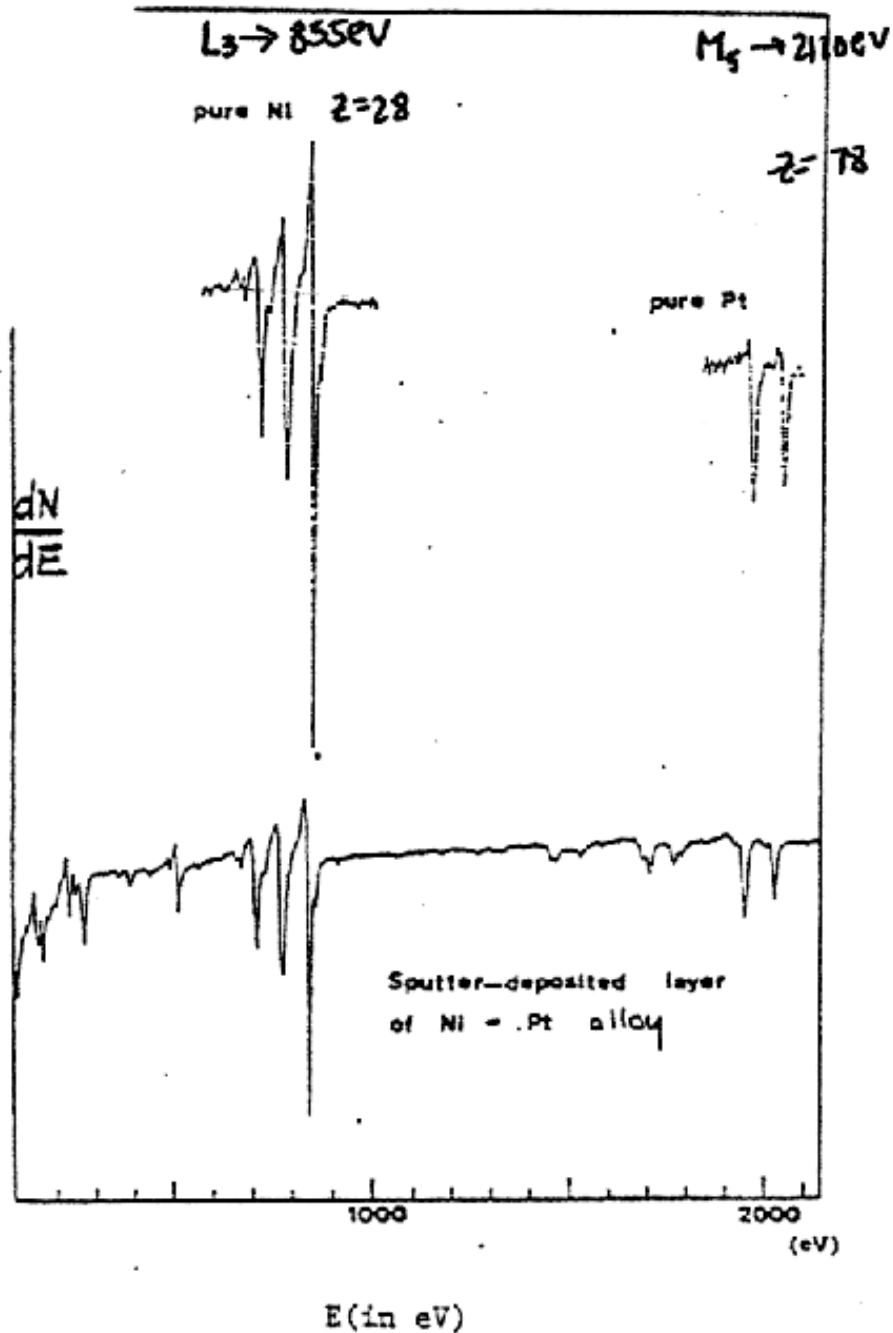


=



LMM

Photo



$$E_0 = 10\text{keV}$$

From R. Shimizu,
Toyota Foundation,
1981

Auger Quantitation . 1.

$$S(WXY) = \underbrace{[n \lambda I_p (1+R^*)]}_{NJ} \underbrace{\sigma(E_w, E_0)}_{\sigma} \underbrace{\delta(WXY)}_{Y} \underbrace{[\Gamma_R^*]}_F$$

We can do perform quantitation in one of three ways:

- 1) using 1st principles - the formula above
- 2) using standards
- 3) using "sensitivity factors"

In order to use (1), we need a good handle on all the terms above: the mean free path of the escaping Auger electrons, the backscattering correction, the inelastic cross-section, the Auger yield and the efficiency factor, F .

for (2): we require the use of "known" standards, where the composition is known.

for (3): we require measurements of Auger signals from a wide variety of known samples that are ~~not~~ normalized to one particular Auger signal.

Anger Quantitation. 2

Our quantitation equation becomes:

$$\boxed{\frac{n_1}{n_2} = \frac{S_1}{S_2} \times \frac{K_2}{K_1}} \quad \text{or} \quad \underline{S_1 = n_1 K_1, S_2 = n_2 K_2}$$

where $K_1 = [\lambda I_p (1+R^*)], \sigma_1 (E_w, E_0) \delta_1 [\bar{r}_2^* T],$

We will look at the 1st principle's calculations last, since that is the case where we need to know all of the factors reasonably well.

CASE 2. The method of Standards

$$\frac{S_A}{S_{A \text{ pure}}} = \frac{n_A K_A}{n_{A \text{ pure}} K_{A \text{ pure}}}$$

where $n = \#/\text{volume}$

Auger Quantitations. 3.

<case 2, cont>

To the extent that the only factors that might be different between K_A and $K_{A \text{ pure}}$ are $(1+R^*)$, λ and $[F_R^* J]$

and if those factors aren't much different, THEN

$$\boxed{S_A / S_{A \text{ pure}} \cong \frac{N_A}{N_{A \text{ pure}}}} \quad \text{with approx.}$$

$$\text{then } \frac{N_A}{N_B} \cong \frac{N_{A \text{ pure}}}{N_{B \text{ pure}}} \frac{S_A}{S_{A \text{ pure}}} \frac{S_{B \text{ pure}}}{S_B}$$

$$\therefore \frac{N_A}{N_B} \cong \left[\frac{N_{A \text{ pure}}}{N_{B \text{ pure}}} \frac{S_{B \text{ pure}}}{S_{A \text{ pure}}} \right] \frac{S_A}{S_B}$$

Better approximation:

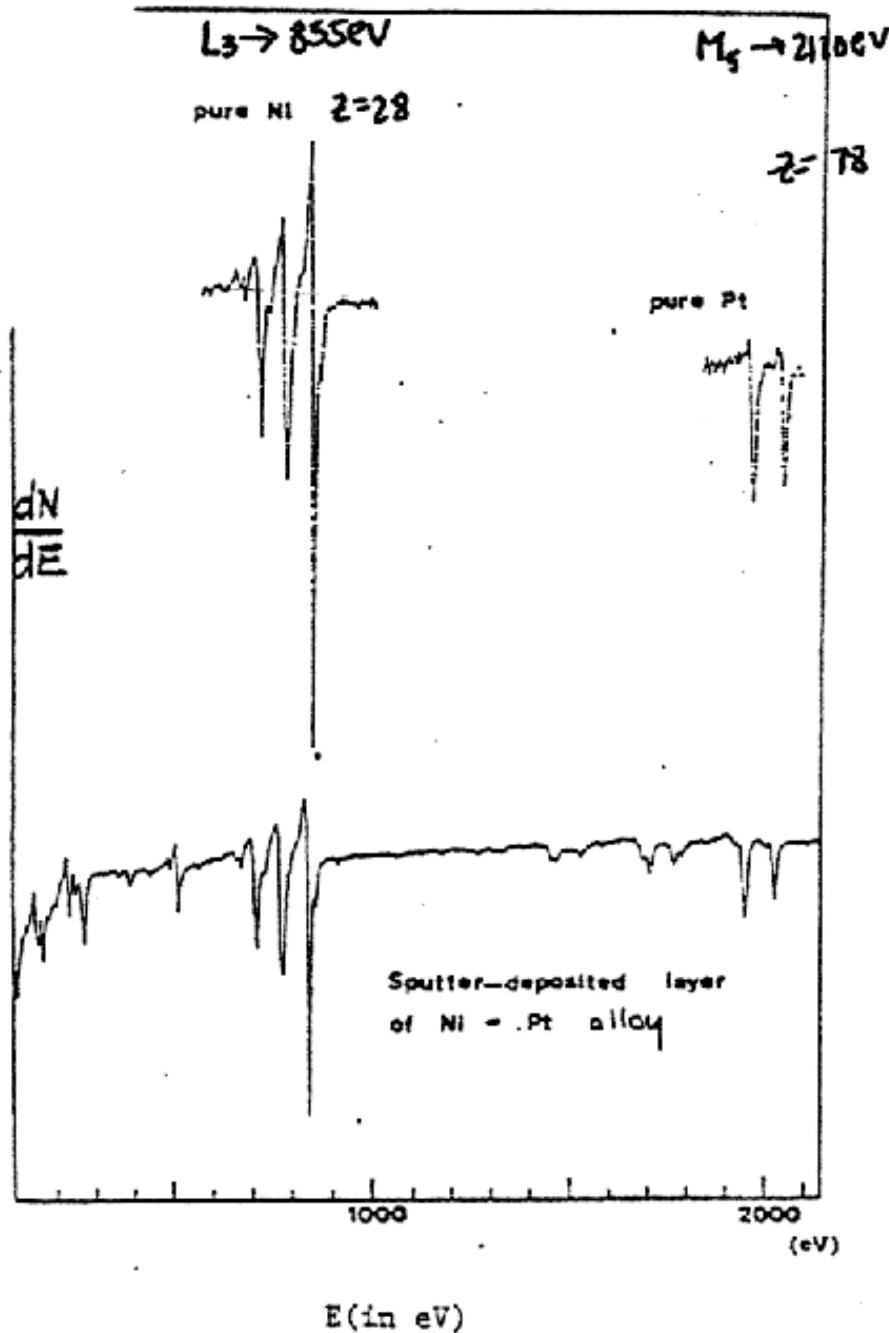
assuming that $\lambda_A \cong \lambda_{\text{pure}}$, not so unreasonable
~~then~~ AND that $[F_R^* J]$ are the same (OK since
the detector is looking at the same energy Auger
electrons and one can prepare samples with similar
topography).

$$\therefore \frac{N_A}{N_B} = \left[\frac{N_{A \text{ pure}}}{N_{B \text{ pure}}} \right] \left[\frac{S_{B \text{ pure}}}{S_{A \text{ pure}}} \right] \left[\frac{S_A}{S_B} \right] \left[\frac{(1+R^*)_{A \text{ pure}}}{(1+R^*)_{B \text{ pure}}} \right] \left[\frac{(1+R^*)_B}{(1+R^*)_A} \right]$$

↑
pure samples

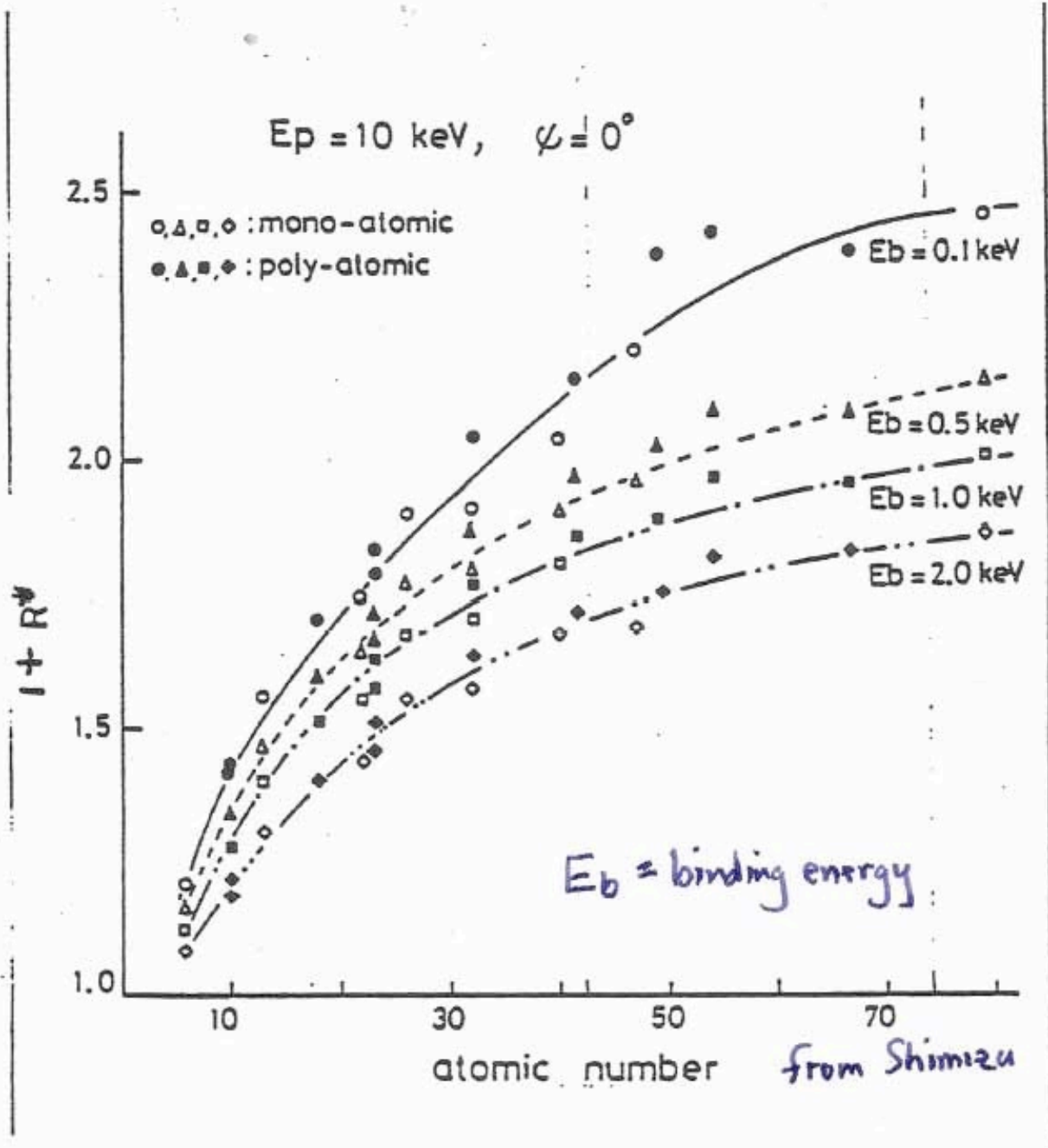
LMM

Photo

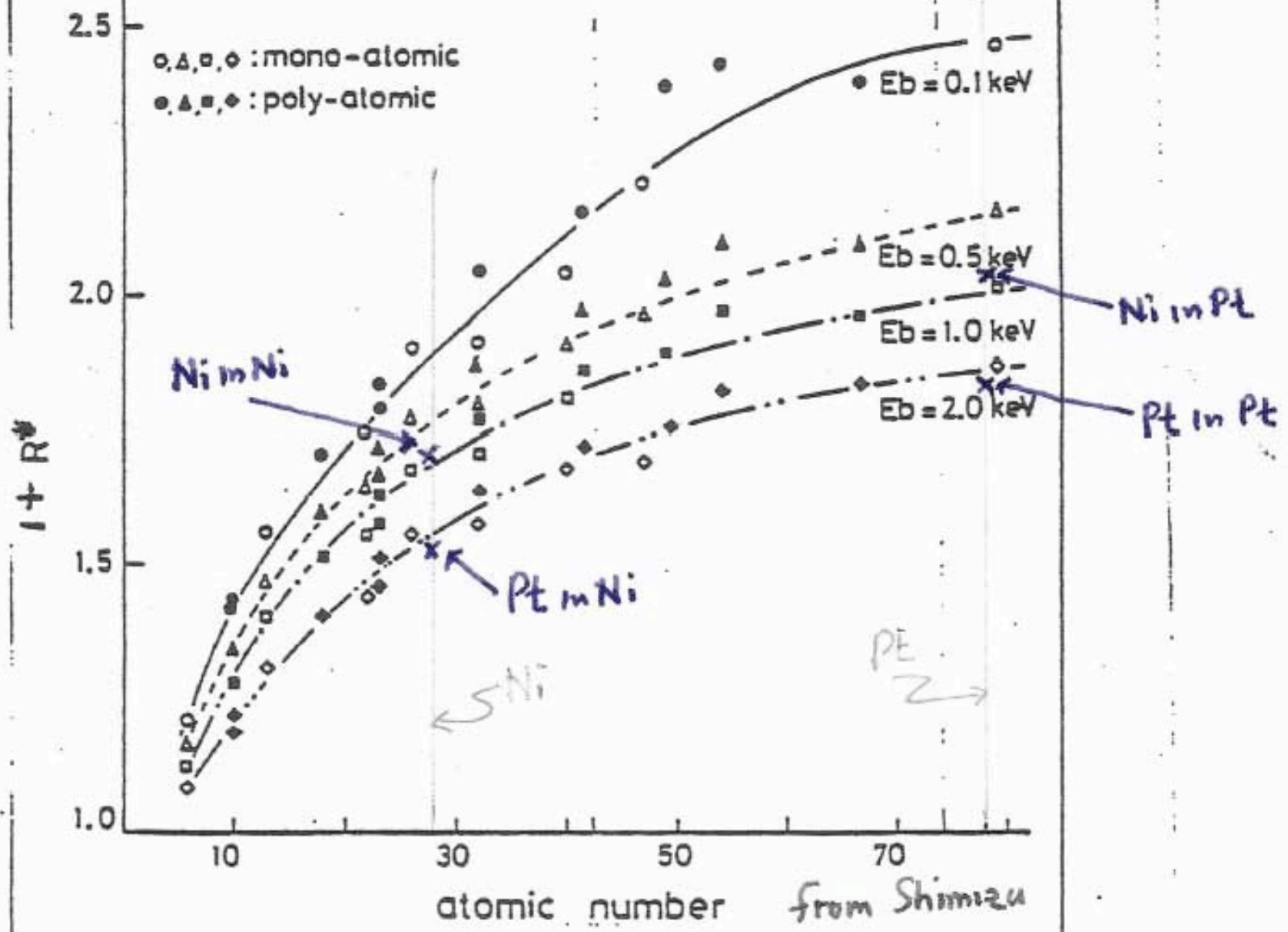


$$E_0 = 10\text{keV}$$

From R. Shimizu,
Toyota Foundation,
1981



$E_p = 10 \text{ keV}, \psi = 10^\circ$



Anger Quantitation 5.

<case 2, cont> method of standards.

so sputter deposited Ni-Pt alloy..

$Z_{Ni} = 28$, $Z_{Pt} = 78$ so R^* correction is needed.

$$\therefore \frac{n_{Ni pms}}{n_{Pt pms}} = \frac{\rho_{Ni pms} \cdot A_{Pt}}{\rho_{Pt pms} \cdot A_{Ni}} \quad , \quad A = \text{atomic weight}$$

$$\therefore \boxed{\frac{n_{Ni pms}}{n_{Pt pms}} = 1.37}$$

So 1st guess is:

$$\frac{n_{Ni}}{n_{Pt}} \cong 1.37 \left[\frac{S_{Pt pms}}{S_{Ni pms}} \right] \left[\frac{S_{Ni}}{S_{Pt}} \right]$$
$$\cong 1.37 \times 0.25 \times 4.5 = \boxed{1.55 = \frac{n_{Ni}}{n_{Pt}}} \leftarrow$$

$$\therefore C_{Ni} \cong \frac{n_{Ni}}{n_{Ni} + n_{Pt}} = \frac{1}{1 + \left(\frac{n_{Pt}}{n_{Ni}}\right)} = \boxed{0.61 = C_{Ni}}$$

we thus to estimate $(1+R^*)$ factor
noting that

$$(1+R^*)_{A \text{ in } [A, B]} = C_A (1+R^*)_{A \text{ in } A} + (1-C_A) (1+R^*)_{A \text{ in } B}$$

Anger Quantitation. 6

we can get the $(1+R^*)$ from the curves by Shimizu
(which are combinations of measurements and
MC calculations/simulations)

$$\begin{aligned} \therefore (1+R^*)_{Ni \text{ in } Ni} &= 1.70 \\ (1+R^*)_{Ni \text{ in } Pt} &= 2.10 \\ (1+R^*)_{Pt \text{ in } Ni} &= 1.51 \\ (1+R^*)_{Pt \text{ in } Pt} &= 1.83 \end{aligned} \left. \begin{array}{l} \\ \\ \\ \end{array} \right\} \begin{array}{l} \\ \\ \end{array} \left. \begin{array}{l} \\ \\ \end{array} \right\} \begin{array}{l} \text{the 2 extremes} \\ \end{array}$$

using our 1st guess at $C_{Ni} = 0.61$ we get

$$(1+R^*)_{Ni \text{ in } NiPt} \cong (0.61)(1.70) + (0.39)(2.10) = 1.86$$

$$(1+R^*)_{Pt \text{ in } NiPt} \cong (0.39)(1.83) + (0.61)(1.51) = 1.63$$

∴ next estimate becomes

$$\begin{aligned} \frac{N_{Ni}}{N_{Pt}} &= (1.55) \left[\frac{(1+R^*)_{Ni \text{ in } Ni}}{(1+R^*)_{Pt \text{ in } Pt}} \right] \left[\frac{(1+R^*)_{Pt \text{ in } NiPt}}{(1+R^*)_{Ni \text{ in } NiPt}} \right] \\ &\quad \uparrow \text{1st guess} \\ &= (1.55) \left[\frac{1.70}{1.83} \right] \left[\frac{1.63}{1.86} \right] = 1.262 \end{aligned}$$

$$\therefore C_{Ni} = \frac{1}{1 + \left(\frac{N_{Pt}}{N_{Ni}}\right)} = \boxed{0.558} = C_{Ni} \quad 10\% \text{ diff than before.}$$

one more iteration gives no more change!
actual concentration, $C_{Ni} = \boxed{0.536}$ 5% accuracy /

Avogadro Quantitation. 7

CASE 3. "relative sensitivity" method

similar to "standards" method except that the k factors are determined by someone ~~else~~ and tabulated.

** care must be taken, since the conditions under which sensitivity factors were measured needs to be similar to conditions of your measurement since it assumes "identical" experimental conditions AND assumes the k_A/k_B are independent of material composition.

generally these "tables" are supplied by manufacturer //

$$\text{From before: } \boxed{\frac{n_A}{n_B} = \frac{S_A k_B}{S_B k_A}} \Rightarrow C_A = \frac{1}{1 + \frac{n_B}{n_A}} = \frac{1}{1 + \frac{S_B k_B}{S_A k_A}}$$

so these "k" factors always appear as a ratio /
∴ we measure the ratio for one standard and everything else is referenced to that.

$$\text{ie, } \underline{\underline{\frac{k_A}{k_C} = \frac{k_A}{k_B} \times \frac{k_B}{k_C}}}$$

Auger Quantitation .7

< case 3, not >

if we tabulate the ratio K_i/K_c for all elements, i then for any ratio,

$$\boxed{\frac{K_i}{K_j} = \left(\frac{K_i}{K_c} \right) / \left(\frac{K_j}{K_c} \right)}$$

the ratio of the
"relative sensitivity"
factors

in the tables (Davis et al being most widely used)

$$\boxed{K_x = K_x / K_{Ag(351eV)}}$$

rel. sensitivity factor — relative to the
silver MNN line at 351eV ///

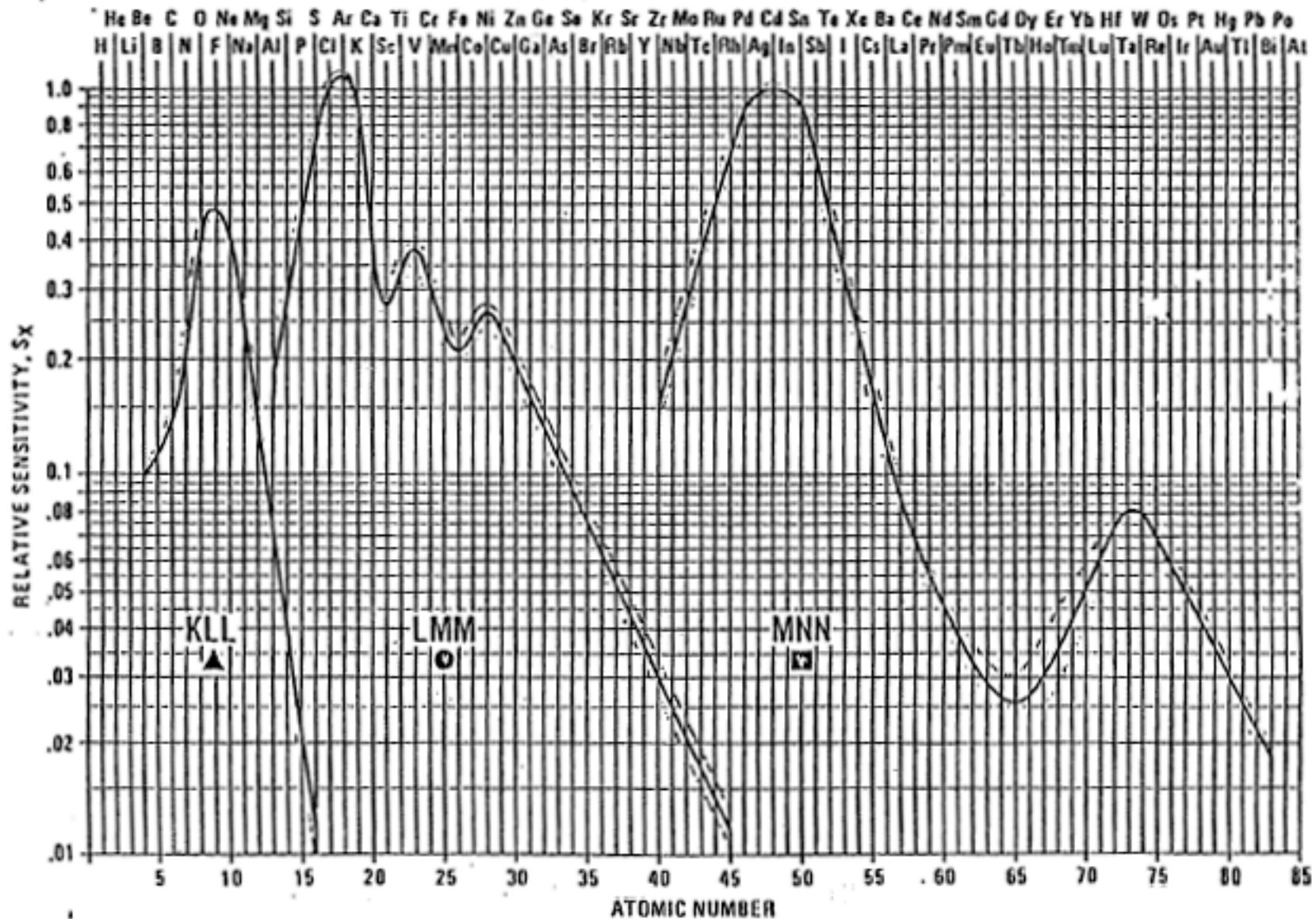
C.E. Davis, et al. Handbook of Auger Electron Spectroscopy. (Physical Electronics Division, Perkin-Elmer Corp, 1978)

similar one / Grant and Briggs. Surface Analysis by Auger and X Ray Photoelectron Spectroscopy (2003). (I.M. Publishers Chichester)

Briggs and Seah, Practical Surface Analysis by Auger and X Ray Photoelectron Spectroscopy (1987) (John Wiley, NY)

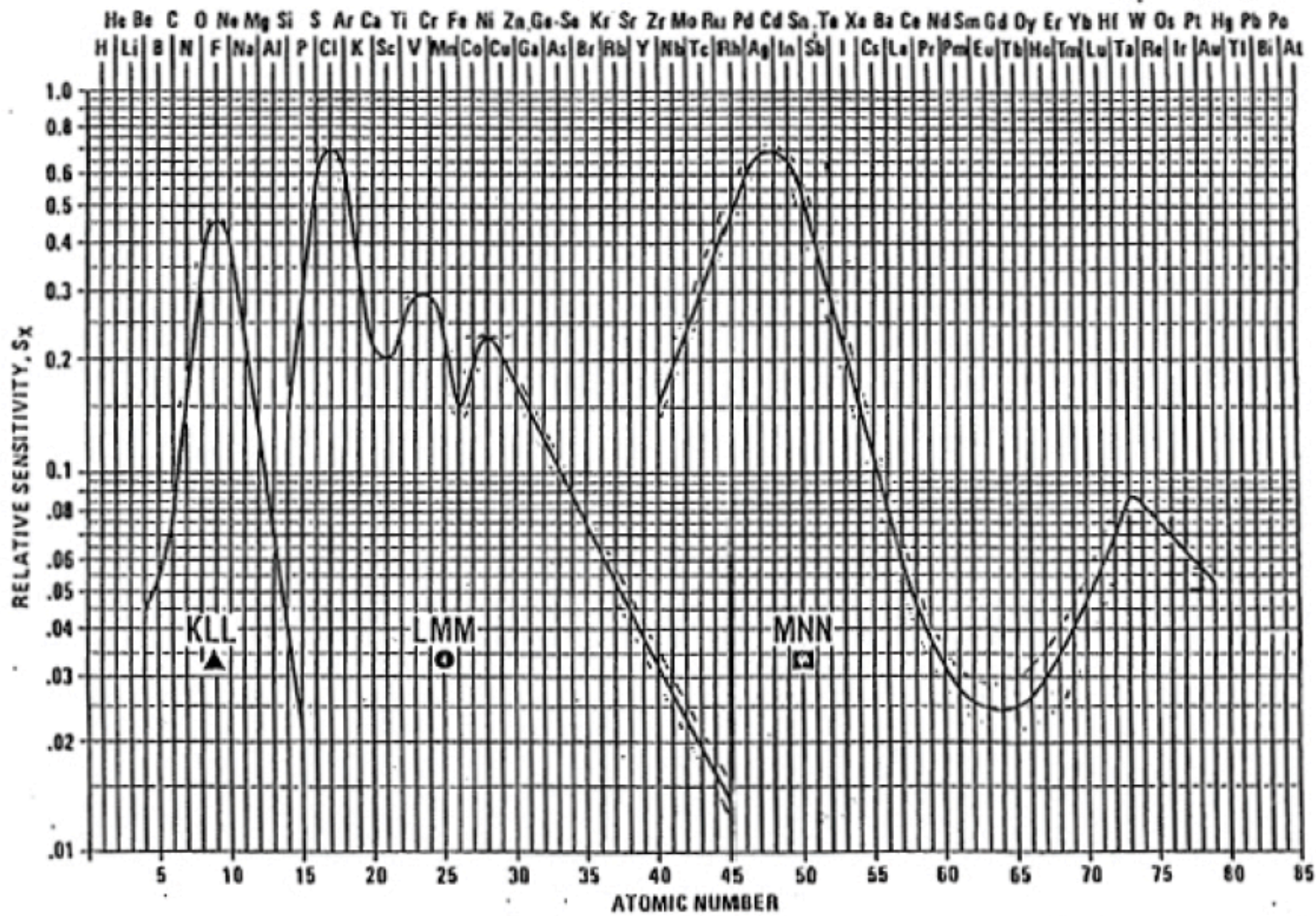
Auger Sensitivity Factors (relative to Ag (MNN), E = 321 eV)

$t_p = 5 \text{ keV}$



Auger Sensitivity Factors (relative to Ag (MNN), $E = 321$ eV)

$E_p = 10$ keV



Auger Quantitation. 8

< case 3 >

let's evaluate the last case, Ni Pt alloy.
using relative sensitivity factors.

$$C_{Ni} = \frac{1}{1 + \left(\frac{S_{Pt}}{S_{Ni}}\right) \left(\frac{K_{Pt}}{K_{Ni}}\right)} = \frac{1}{1 + (4.5) \left(\frac{K_{Pt}}{K_{Ni}}\right)}$$

from Davis tables (curves):

$$\frac{K_{Ni}}{K_{Pt}} = \frac{K_{Ni}}{K_{Ag}} \left(\frac{K_{Pt}}{K_{Ag}} \right) = \frac{.225}{.0575} = 3.91 = \frac{K_{Ni}}{K_{Pt}}$$

$$\therefore C_{Ni} = \frac{1}{1 + (4.5)(3.91)} = 0.54 = C_{Ni}$$

we can extend this to multielement systems:

if n_i = # density of the i^{th} component,
the atomic concentration is then

$$C_i = \frac{n_i}{\sum_j n_j} = \frac{S_i/K_i}{\sum_j (S_j/K_j)} = \frac{1}{\sum_j \left(\frac{S_j}{S_i}\right) \left(\frac{K_i}{K_j}\right)} = C_i$$

NOTE: the relative sensitivity factors don't necessarily work for the different backscattering conditions. also, they are dependent upon the incident electron energy.